

Polynomial time approximation of entropy of shifts of finite type

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- **Full shift:** $\mathcal{A}^{\mathbb{Z}^d}$ over a finite alphabet \mathcal{A} .
- **Shift space:** for some list \mathcal{F} of “forbidden” configurations on finite shapes,
 $X = X_{\mathcal{F}} := \{x \in \mathcal{A}^{\mathbb{Z}^d} : x \text{ contains no elements of } \mathcal{F}\}$
- **Shift of finite type (SFT):** a shift space where \mathcal{F} can be chosen finite.
- **Nearest neighbor (n.n.) SFT:** a shift space where all elements of \mathcal{F} are configurations on *edges* of \mathbb{Z}^d .

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Topological entropy

- $B_n := [0, n-1]^d$

- **globally admissible** configurations:

$$GA_n(X) = \{x(B_n) : x \in X\}$$

- **Topological entropy:**

$$h(X) := \lim_{n \rightarrow \infty} \frac{\log |GA_n(X)|}{n^d}$$

- **locally admissible** configurations:

$$LA_n(X) = \{ \text{configs. on } B_n \text{ forbidding } \mathcal{F} \}$$

- Theorem (Ruelle, Friedland):

$$h(X) = \lim_{n \rightarrow \infty} \frac{\log |LA_n(X)|}{n^d}$$

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- A \mathbb{Z} n.n. SFT X over alphabet \mathcal{A} is specified by a directed graph G with vertices indexed by \mathcal{A} and an edge from a to b iff $ab \notin \mathcal{F}$.

Golden Mean Shift:

- Adjacency matrix A of G is the square matrix indexed by \mathcal{A} :

$$A_{ab} = \begin{cases} 1 & ab \notin \mathcal{F} \\ 0 & ab \in \mathcal{F} \end{cases}$$

- $h(X) = \log \lambda(A)$, where $\lambda(A)$ is the spectral radius of A .
- Characterization of entropies for $d = 1$ (Lind):

$$\{\log \lambda^{1/q}\}$$

where λ is a Perron number and $q \in \mathbb{N}$

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Examples of \mathbb{Z}^2 SFTs: hard squares

- **hard squares** $\mathcal{A} = \{0, 1\}, \mathcal{F} = \left\{ \begin{array}{c} 11 \\ 1 \end{array} \right\}$
- $h(\text{hard squares}) = ???$
- $h(\text{hard hexagons}) = \log(\lambda)$ where λ is an algebraic integer of degree 24.

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Examples of \mathbb{Z}^2 SFTs: checkerboard (coloring) constraints

- **q -checkerboard \mathcal{C}_q :** $\mathcal{A} = \{1, \dots, q\}, \mathcal{F} = \{aa, \begin{smallmatrix} a \\ a \end{smallmatrix}\}$
- $h(\mathcal{C}_2) = 0$
- (Lieb): $h(\mathcal{C}_3) = (3/2) \log(4/3)$
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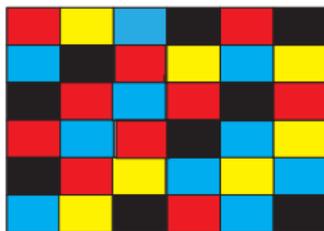
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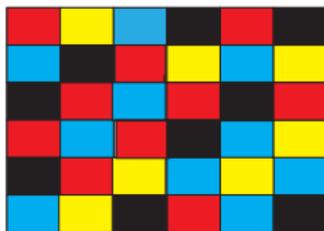
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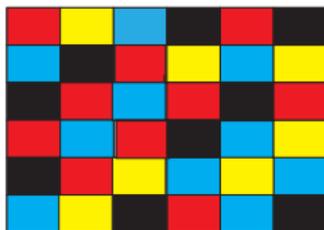
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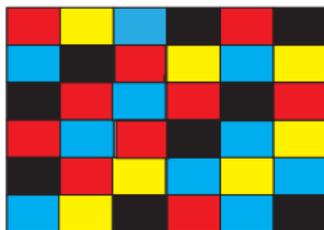
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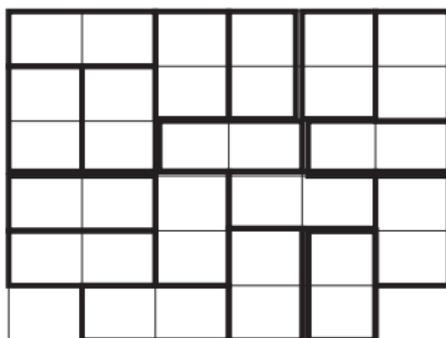
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Examples of \mathbb{Z}^2 SFT's: dimers

- **dimers:**



$$\mathcal{F} = \{LL, LT, LB, RR, TR, BR, \begin{matrix} T & T & T & B & L & R \\ L & R & T & B & B & B \end{matrix}\}$$

- (Fisher-Kastelyn-Temperley (1961)):

$$h(\text{Dimers}) = \frac{1}{16\pi^2} \int_{-\pi}^{\pi} \int_{-\pi}^{\pi} \log(4 + 2\cos\theta + 2\cos\phi) d\theta d\phi$$

- $h(\text{Monomers-Dimers}) = ???$

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L	R	T	T	T	T
T	T	B	B	B	B
B	B	L	R	L	R
L	R	T	L	R	T
L	R	B	T	T	B
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Ledrappier 3-dot

$$X = \{x \in \{0, 1\}^{\mathbb{Z}^2} : x((i, j)) + x((i+1, j)) + x((i, j+1)) = 0 \pmod{2}\}$$

$$\mathcal{F} = \left\{ \begin{array}{c} a \\ b \\ c \end{array} : a + b + c \neq 0 \pmod{2} \right\}$$

$$h(X) = 0.$$

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Examples of \mathbb{Z}^2 SFTs: iceberg model

- $\mathcal{A} = \{-M, \dots, -1, 0, 1, \dots, M\}$
- $\mathcal{F} = \{ab, \begin{smallmatrix} a \\ b \end{smallmatrix} : a, b \text{ have opposite signs}\}$
- positives can sit next to positives and zeros, negatives can sit next to negatives and zeros, and zeros can sit next to anyone.
- Example: $M = 2$

1	2	2	1	0	0	-1	-2	0	2
0	2	1	0	-1	0	-2	-1	0	2
1	1	0	-2	0	0	0	0	1	2
0	2	1	0	-1	0	-2	-1	0	2
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Topological entropy, $d \geq 2$

- Exact formula known only in a few cases.
- Characterization of entropies for $d \geq 2$ (Hochman-Meyerovitch):

{**right recursively enumerable** (RRE) numbers $h \geq 0$ }

i.e, there is an algorithm that produces a sequence $r_n \geq h$
s.t. $r_n \rightarrow h$.

Proof:

- Necessity: Let $r_n := \frac{\log |LA_n|}{n^d}$.
By Ruelle/Friedland Theorem, $r_n \rightarrow h$.
By subadditivity of $\log |LA_n|$, each $r_n \geq h$.
 - Sufficiency (hard): Emulate Turing machine with an SFT.
- RRE's can be poorly computable, or even non-computable.

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- Exact formula known only in a few cases.
- Characterization of entropies for $d \geq 2$ (Hochman-Meyerovitch):

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- $h(\mu) := \lim_{n \rightarrow \infty} \frac{H_\mu(B_n)}{n^d}$
- $d = 1$: Theorem: $h(\mu) = H_\mu(0 \mid \{-1, -2, -3, \dots\})$
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For $\bar{z} \in \mathbb{Z}^2$, let $\mathcal{P}^-(\bar{z}) := \{\bar{z}' \in \mathbb{Z}^2 : \bar{z}' \prec \bar{z}\}$ the lexicographic past of \bar{z} , and $\mathcal{P}^- := \mathcal{P}^-(0)$



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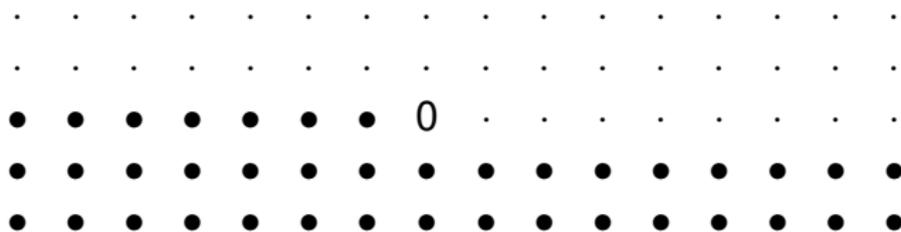


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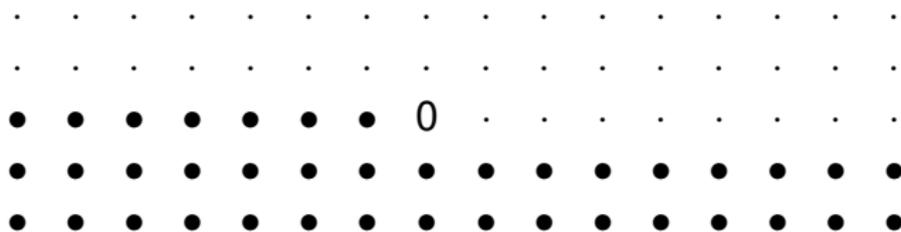


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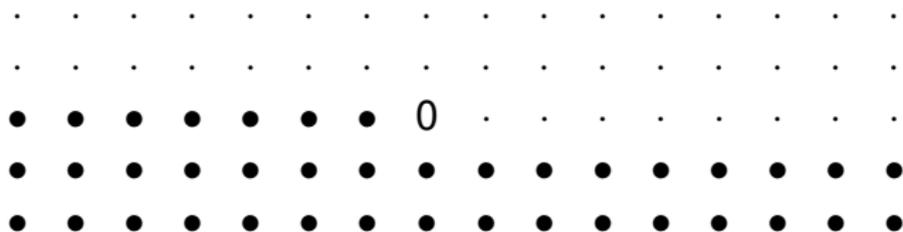
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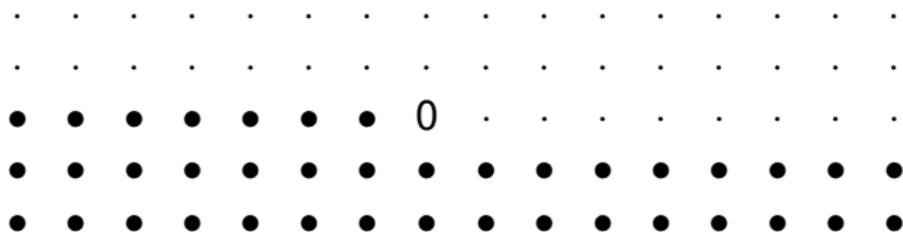
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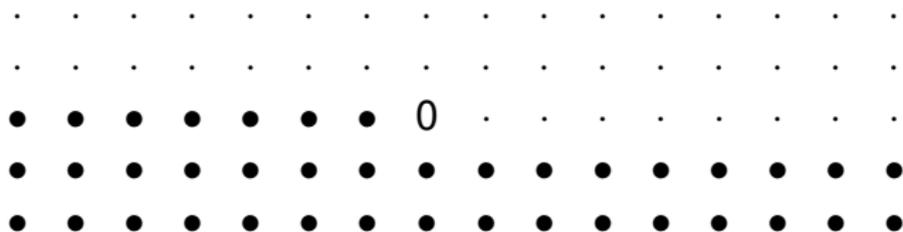
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Variational Principle for Topological Entropy

- For a shift space X ,

$$h(X) = \sup_{\mu} h(\mu)$$

where the sup is taken over all shift-invariant Borel probability measures μ s.t. $\text{support}(\mu) \subseteq X$.

- Fact: The sup is always achieved. A measure which achieves the sup is called a **measure of maximal entropy (MME)**.
- So for an MME μ , $h(X) = h(\mu) = \int I_{\mu}(x) d\mu(x)$
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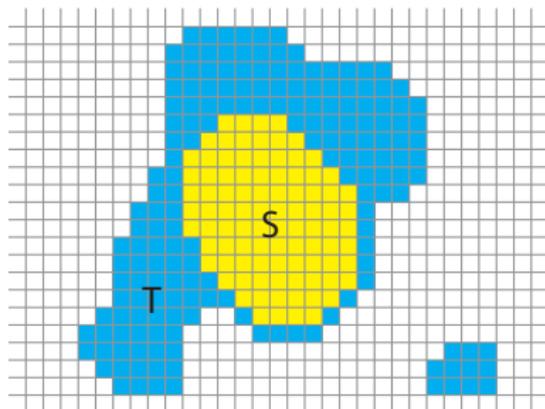
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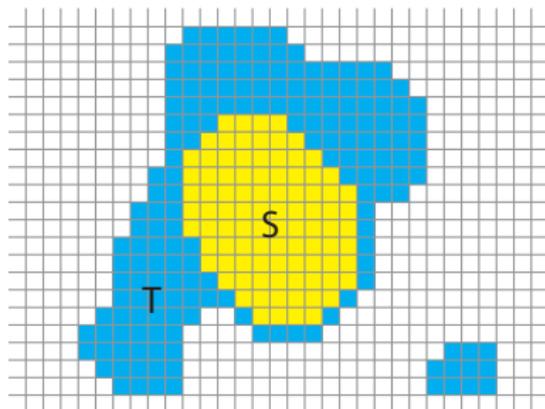
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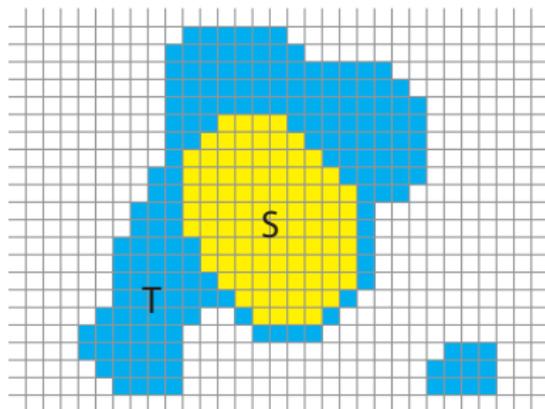
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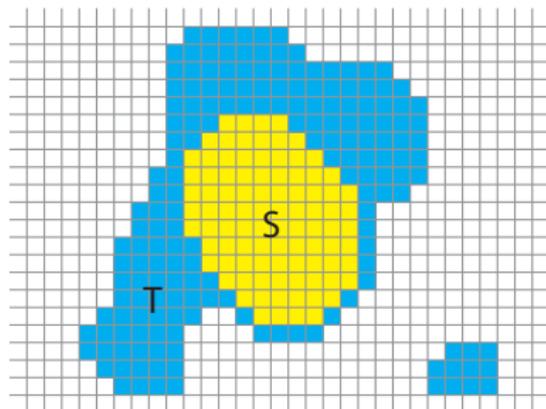
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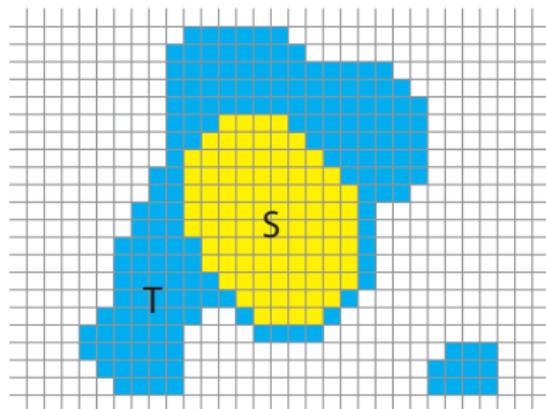
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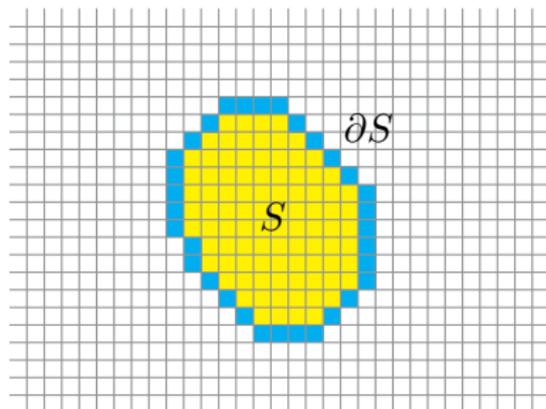
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Theorem (Lanford/Ruelle, Burton/Steif): Every MME on a n.n. SFT is a uniform MRF.

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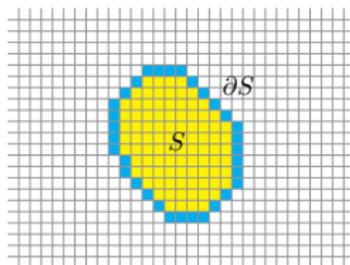
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Theorem (Lanford/Ruelle, Burton/Steif): Every MME on a n.n. SFT is a uniform MRF.

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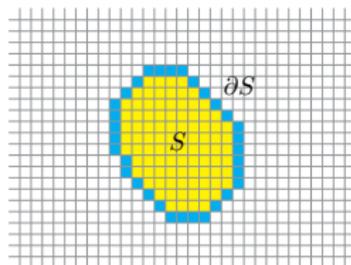
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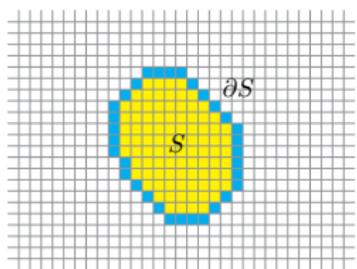
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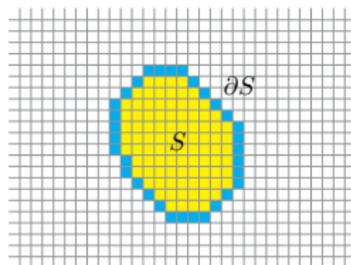
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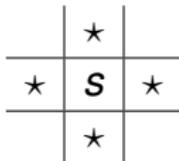


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Entropy Representation

Let $R_{a,b,c} := [-a, -1] \times [1, c] \cup [0, b] \times [0, c]$

Example: $R_{3,4,3}$:



Theorem (special case of Gamarnik-Katz): Let X be a n.n. \mathbb{Z}^d SFT and μ an MME on X . If

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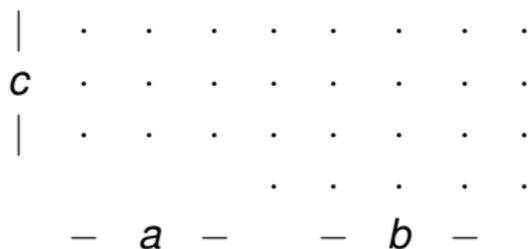
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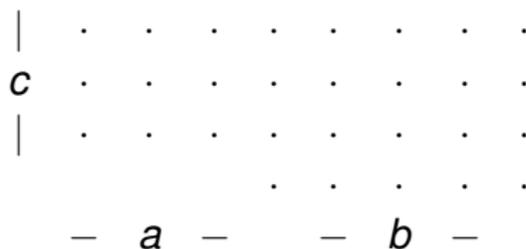
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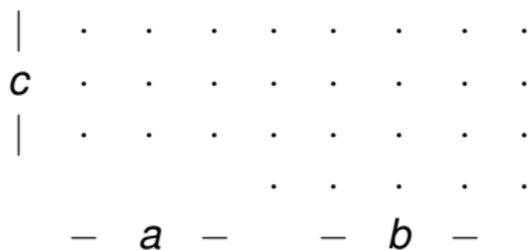
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$$h(X) = \lim_{n \rightarrow \infty} \frac{-\log \mu(s^{B_n} \mid s^{\partial B_n})}{n^d}$$

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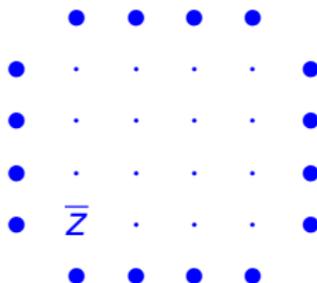
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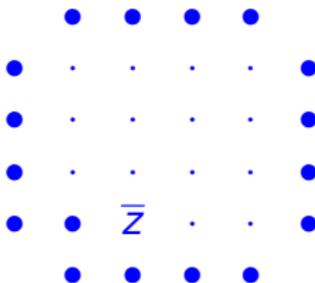
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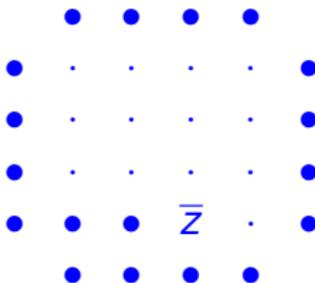
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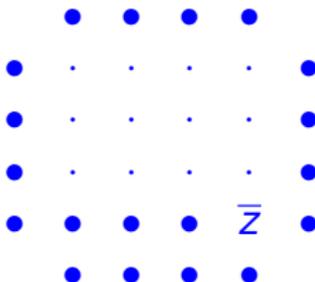
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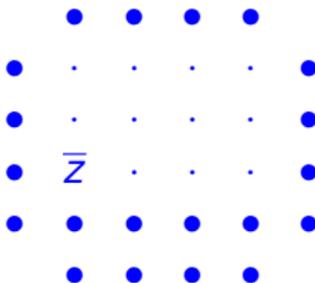
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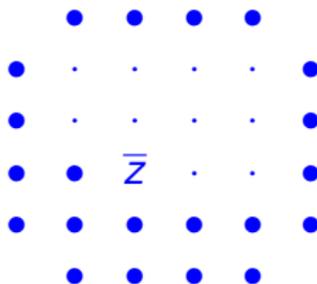
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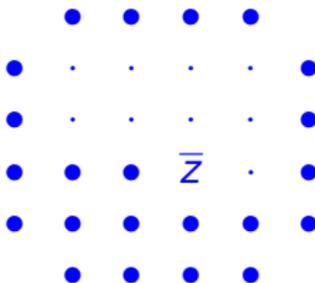
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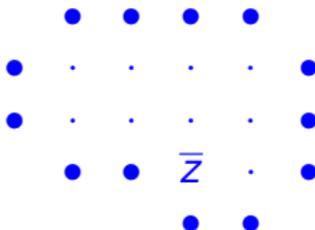
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Algorithmic consequence

Theorem (special case of Gamarnik-Katz): Let X be a n.n. \mathbb{Z}^2 SFT and μ an MME on X . If

- 1 X has a safe symbol s – and –
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$$L := \lim_{a,b,c \rightarrow \infty} \mu(s^0 \mid s^{\partial R_{a,b,c}}) \text{ exists}$$

and convergence is exponential

Then there is a polynomial time algorithm to compute $h(X) = -\log L$.

Proof: Approximate L by $\mu(s^0 \mid s^{\partial R_{n,n,n}})$.

- Accuracy is $e^{-\Omega(n)}$
- Claim: Computation time is $e^{O(n)}$
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Proof: Approximate L by $\mu(s^0 \mid s^{\partial R_{n,n,n}})$.

- Accuracy is $e^{-\Omega(n)}$
- Claim: Computation time is $e^{O(n)}$
- Trade exponential accuracy in exponential time for linear accuracy ($1/n$) in polynomial time. \square

Algorithmic consequence

Theorem (special case of Gamarnik-Katz): Let X be a n.n. \mathbb{Z}^2 SFT and μ an MME on X . If

- 1 X has a safe symbol s – and –
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Proof of Claim, via transition matrices

$$\mu(s^0 \mid s^{\partial R_{n,n,n}}) = \frac{\begin{array}{cccccc} & s & s & s & s & s \\ s & \cdot & \cdot & \cdot & \cdot & \cdot & s \\ \# & s & \cdot & \cdot & \cdot & \cdot & s \\ s & s & s & s & \cdot & \cdot & s \\ & & & s & s & s \\ s & s & s & s & s & s \\ s & \cdot & \cdot & \cdot & \cdot & \cdot & s \\ \# & s & \cdot & \cdot & \cdot & \cdot & s \\ s & s & s & \cdot & \cdot & \cdot & s \\ & & & s & s & s \end{array}}{\begin{array}{cccccc} & s & s & s & s & s \\ s & \cdot & \cdot & \cdot & \cdot & \cdot & s \\ \# & s & \cdot & \cdot & \cdot & \cdot & s \\ s & s & s & \cdot & \cdot & \cdot & s \\ & & & s & s & s \end{array}}$$

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M_i is transition matrix from column i to column $i + 1$ compatible with $s^{\partial S_{n,n,n}}$ and

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Topological Strong Spatial Mixing (TSSM)

- A \mathbb{Z}^d SFT X satisfies **topological strong spatial mixing (TSSM)** with gap g if

for any disjoint $U, S, V \in \mathbb{Z}^d$ s.t. $d(U, V) \geq g$,

$u \in A^U, s \in A^S, v \in A^V$, s.t. us and sv are globally admissible,

then so is usv .

- Safe symbol \Rightarrow TSSM

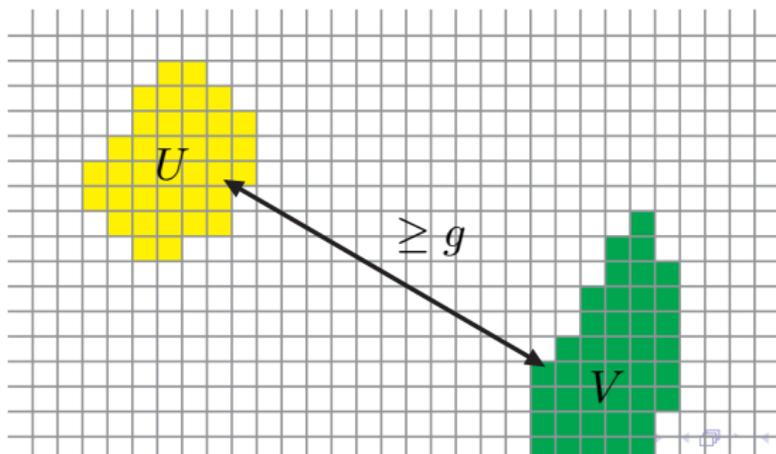
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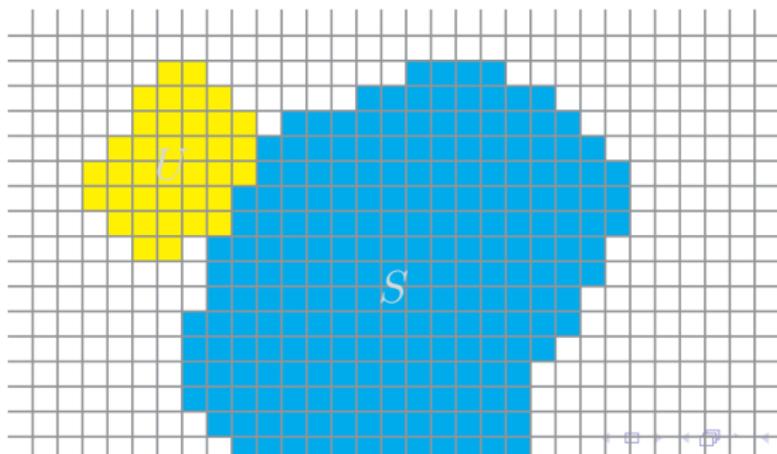
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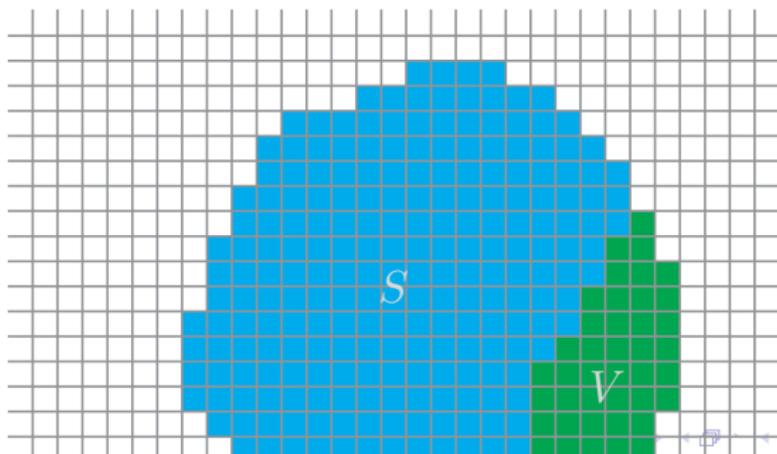
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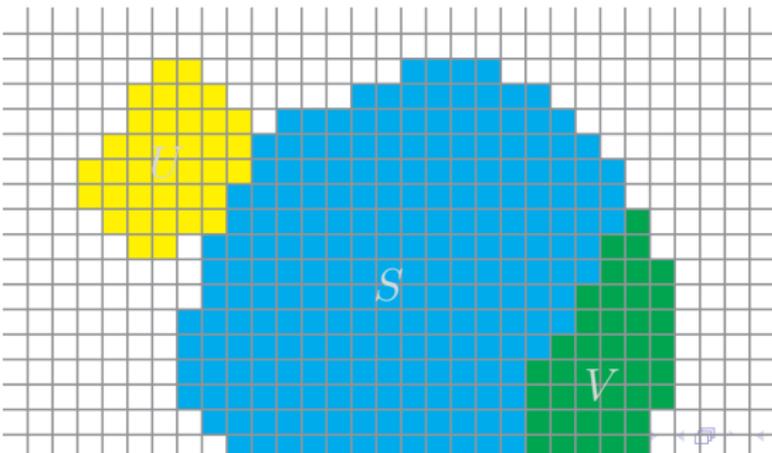
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Generalization

Theorem (Adams, Briceno, Marcus, Pavlov): Let X be a \mathbb{Z}^d n.n. SFT and μ an MME on X . If

- 1 X satisfies TSSM
- 2 For some periodic orbit O in X and all $\omega \in O$

$$L(\omega) := \lim_{a,b,c \rightarrow \infty} \mu(\omega(0) \mid \omega(\partial R_{a,b,c})) \text{ exists}$$

Then

$$h(X) = -\frac{1}{|O|} \sum_{\omega \in O} \log L(\omega)$$

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Verification of exponential convergence condition: using coupling and Peierls arguments.

Applies to:

- hard squares
- q -checkerboard with $q \geq 6$
- iceberg with $M \geq 24$.

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- An MRF μ satisfies **strong spatial mixing (SSM)** at rate $f(n)$

if for all $V \in Z^d$, $U \subset V$

all $u \in A^U$, and $v, v' \in A^{\partial V}$ satisfying $\mu(v), \mu(v') > 0$,

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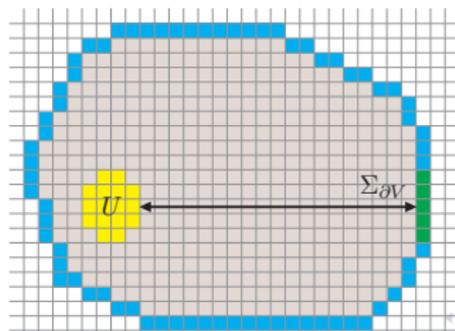
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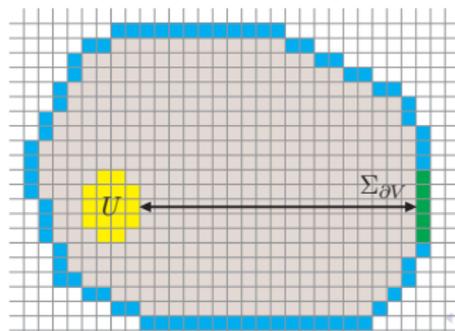
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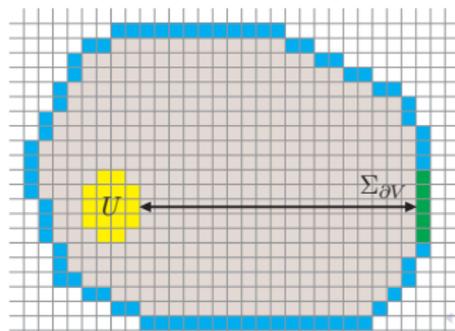
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Stronger conclusion

Theorem (Briceno): Let X be a \mathbb{Z}^d n.n. SFT and μ an MME on X . If

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Then for *all* invariant measures ν s.t. $\text{support}(\nu) \subseteq X$,

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Extension to Pressure

- Generalize results from entropy to *pressure* of nearest neighbour interactions
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Topological Pressure and Variational Principle

- Let X be a shift space and $f : X \rightarrow \mathbb{R}$ a continuous function.
- **Topological Pressure** (defined by Variational Principle):

$$P_X(f) := \sup_{\mu} h(\mu) + \int f d\mu$$

where the sup is taken over all shift-invariant Borel probability measures μ such that $\text{support}(\mu) \subseteq X$.

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- For a nearest-neighbor interaction Φ , the *underlying SFT*:
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- A *nearest neighbour (n.n.) Gibbs measure* μ corresponding to Φ is an MRF on X such that for $S \Subset \mathbb{Z}^d$, $\delta \in \mathcal{A}^{\partial S}$, $\mu(\delta) > 0$, $w \in \mathcal{A}^S$:

$$\mu(w|\delta) = \frac{e^{-U^\Phi(w\delta)}}{Z^{\Phi, \delta}(S)}.$$

where

- $U^\Phi(w\delta)$ is the sum of all Φ -values of $w\delta$ for vertices, edges in $S \cup \partial S$
- $Z^{\Phi, \delta}(S)$ is the normalization factor.

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- A *nearest neighbour (n.n.) Gibbs measure* μ corresponding to Φ is an MRF on X such that for $S \Subset \mathbb{Z}^d$, $\delta \in \mathcal{A}^{\partial S}$, $\mu(\delta) > 0$, $w \in \mathcal{A}^S$:

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Examples of n.n. Gibbs measures

- uniform MME on n.n. SFT
- hard square model with activities
- ferromagnetic Ising model with no external field.

Equilibrium states versus n.n. Gibbs measures

- Pressure of n.n. interaction Φ :

$$P(\Phi) := \lim_{n \rightarrow \infty} \frac{\log Z^\Phi(B_n)}{n^d}$$

where $Z^\Phi(B_n)$ is the “free boundary” normalization.

- Let $A_\Phi(x) := -\Phi(x(0)) - \sum_{i=1}^d \Phi(x(0), x(e_i))$.
- Fact: $P_{X_\Phi}(A_\Phi) = P(\Phi)$.
- Lanford-Ruelle Theorem: Every equilibrium state for A_Φ is a Gibbs measure for Φ .
- Dobrushin Theorem: If X_Φ is strongly irreducible, then every Gibbs measure for Φ is an equilibrium state for A_Φ .
- These theorems hold in much greater generality.

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Pressure representation and approximation

Theorem (Adams, Briceno, Marcus, Pavlov): Let μ a Gibbs measure for a n.n. interaction Φ with underlying \mathbb{Z}^d n.n. SFT X .
If

- 1 X satisfies TSSM
- 2 For some periodic orbit O in X and all $\omega \in O$

$$L(\omega) := \lim_{a,b,c \rightarrow \infty} \mu(\omega(0) \mid \omega(\partial R_{a,b,c})) \text{ exists}$$

Then

$$P(\Phi) = \frac{1}{|O|} \sum_{\omega \in O} -\log L(\omega) + A_{\Phi}(\omega)$$

Moreover, if $d = 2$ and convergence in hypothesis 2 is exponential, then there is a polynomial time algorithm to compute $P(\Phi)$.

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Stronger conclusion

Theorem (Briceno): Let μ a Gibbs measure for a n.n. interaction Φ with underlying \mathbb{Z}^d n.n. SFT X . If

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Then for all shift-invariant measures ν such that $\text{support}(\nu) \subseteq X$,

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An SFT X satisfies the **D-condition** if

- there exist sequences of finite subsets $(\Lambda_n), (M_n)$ of \mathbb{Z}^d such that $\Lambda_n \nearrow \infty$, $\Lambda_n \subseteq M_n$, $\frac{|M_n|}{|\Lambda_n|} \rightarrow 1$, such that
- for any globally admissible $v \in \mathcal{A}^{\Lambda_n}$ and finite $S \subset M_n^c$ and globally admissible $w \in \mathcal{A}^S$, we have that vw is globally admissible.

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Connection with Thermodynamic Formalism

Theorem: Let μ a Gibbs measure for a n.n. interaction Φ with underlying \mathbb{Z}^d n.n. SFT X . If

- X satisfies the D-condition
- $I_\mu = A_\Psi$ for some *absolutely summable* interaction Ψ s.t.
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MME, $d = 1$

- Assuming adjacency matrix A is irreducible and aperiodic, there is a unique MME μ_{\max} , which is a Markov chain given by transition matrix

$$P_{ij} = \begin{cases} \frac{r_j}{\lambda r_i} & ij \notin \mathcal{F} \\ 0 & ij \in \mathcal{F} \end{cases}$$

where $\lambda = \lambda(A)$ and r is a right eigenvector for λ , and stationary vector $r_i \ell_j$ where ℓ is a left eigenvector for λ (suitably normalized)

- Thus, if $\mu(w_1 w_2 \dots w_{n-1} w_n) > 0$, then

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Entropy representation for MME, $d = 1$



$$\begin{aligned}I_{\mu}(x) &= -\log \mu(x(0) | x(\mathcal{P}^-)) \\ &= -\log P_{x_0 x_{-1}} \\ &= \log \lambda + \log r_{x_{-1}} - \log r_{x_0}\end{aligned}$$

- So, for *all* invariant measures ν ,

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In particular, if the SFT has a fixed point $x^* := a^{\mathbb{Z}}$ and ν is the delta measure on x^* , then on

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- So, for *all* invariant measures ν ,

$$\begin{aligned}\int I_{\mu}(x) d\nu(x) &= \int (\log \lambda + \log r_{x_{-1}} - \log r_{x_0}) d\nu(x) \\ &= \log \lambda \\ &= h(X)\end{aligned}$$

In particular, if the SFT has a fixed point $x^* := a^{\mathbb{Z}}$ and ν is the delta measure on x^* , then on

$$h(X) = \int I_{\mu}(x) d\nu(x) = I_{\mu}(x^*) = -\log \mu(x^*)$$

and so $h(X)$ can be computed from the value of the information function at only one point.

- In this case, $I_{\mu}(x)$ is defined everywhere.